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Global-Mode Analysis of Full-Disk Data from the Michelson Doppler Imager and the Helioseismic and Magnetic Imager

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Abstract Building upon our previous work, in which we analyzed smoothed and subsampled velocity data from the Michelson Doppler Imager (MDI), we extend our analysis to unsmoothed, full-resolution MDI data. We also present results from the Helioseismic and Magnetic Imager (HMI), in both full resolution full resolution and processed to be a proxy for the low-resolution MDI data. We find that the systematic errors we saw previously, namely peaks in both the high-latitude rotation rate and the normalized residuals of odd a-coefficients, are almost entirely absent in both the two full-resolution analyses. Furthermore, we find that both systematic errors seem to depend almost entirely on how the input images are apodized, rather than on resolution or smoothing. Using the full-resolution HMI data, we confirm our previous findings regarding the effect of using asymmetric profiles on mode parameters, and also find that they occasionally result in more stable fits. We also confirm our previous findings regarding discrepancies between 360-day and 72-day analyses. We further investigate a six-month period previously seen in f-mode frequency shifts using the low-resolution datasets, this time accounting for solar-cycle dependence using magnetic-field data. Both the HMI and MDI saw prominent six-month signals in the frequency shifts, but we were surprised to discover that the strongest signal at that frequency occurred in the mode coverage for the low-resolution proxy. Finally, a comparison of mode parameters from the HMI and MDI shows that the frequencies and a-coefficients agree closely, encouraging the concatenation of the two datasets.

Keywords: Helioseismology, Observations; Oscillations, Solar

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#### 1. Introduction

Designed to be the successor to the Michelson Doppler Imager (MDI: Scherrer et al., 1995) onboard the Solar and Heliospheric Observatory (SOHO), the Helioseismic and Magnetic Imager (HMI: Schou et al., 2012) was launched onboard the Solar Dynamics Observatory (SDO) in February 2010. The designs of the two instruments are quite similar; here we shall note the differences between the two projects that are most pertinent to global-mode analysis. The HMI is equipped with a  $4096 \times 4096$  pixels pixel CCD and takes images with a spatial resolution of approximately 0.5 arscec per pixel, or about four times that of MDI. the MDI. The SDO is in geosynchronous orbit, whereas SOHO orbits the Sun–Earth L<sub>1</sub> Lagrange point; partially partly for this reason, the HMI is able to send down much more telemetry. Among other observables, the HMI produces full-resolution dopplergrams at a cadence of 45 seconds. Lastly, Last, the HMI observes the Fe I 6173 Å spectral line, so that it sees a slightly lower height in the solar atmosphere than the MDI, which observed the Ni I 6768 Å line (Fleck, Couvidat, and Straus, 2011).

Global-mode analysis of data from  $\overline{\text{MDI}}$ 's-the  $\overline{\text{MDI}}$  Medium- $\ell$  Program and systematic errors therein were described by Larson and Schou (2015), hereafter referred to as LS15. Before an attempt is made to extend this medium- $\ell$  analysis to HMI data, it is fitting to apply it to the MDI full-disk data and compare the results. Although one might expect the two MDI analyses to be in near-perfect agreement, our investigation reveals surprising differences. In particular, systematic errors such as the "bump" seen in the normalized residuals of the odd a-coefficients and the anomalous peak in the near-surface rotation rate at high latitudes have different characteristics in the analysis of full-disk data.

MDI full-disk data are available throughout the mission, but usually with a low duty cycle. Nominally, for two months per year telemetry was allocated to send down the full-disk images continuously. These time intervals constitute the Dynamics Program. As discussed in the next section, the actual lengths of the full-disk observing campaigns varied widely across the mission, as well as did their timing within the year.

One might say that the primary difference between the MDI full-disk data and low-resolution data (labelled\_labeled\_vw\_V, see LS15) is that the latter are smoothed and subsampled (see Section 3), leaving them with a resolution one-fifth that of the full-disk data. However, another important difference lies in the data cropping. Whereas the vw\_V data are cropped to 90 % of the average solar radius onboard the spacecraft, the full-disk data extend significantly closer to the limb. Further details are provided in Section 3.

In order to provide continuity with the MDI Medium- $\ell$  Program, we use used the HMI data to create a vw\_V proxy. This also allows allowed us to further investigate periodicities seen in the f-mode frequencies from the analysis of MDI vw\_V data.

In the next section we describe the datasets used in our analysis. In Section 3 we discuss how these data were analyzed, with emphasis on how each analysis differs from the analysis in LS15. Section 4 gives the results, first for the MDI and then for the HMI, followed by a comparison of the two instruments. Section

5 describes a six-month periodicity in data from both the MDI and HMI and discusses the effect of  $B_0$  (the heliographic latitude of the sub-observer point) on leakage matrices and the resulting inversions for solar rotation. Finally, in Section 6 we discuss our findings and propose how we might move forward proceed.

#### 2. Data

Beginning in 1996, the MDI was continuously operated in full-disk mode for a few months each year through 2010. We therefore have 15 time intervals to analyze, known as the dynamics runs. To choose the exact intervals to use for global-mode analysis, one must balance the lengths of the timeseries time series and their duty cycles. For the most part, we have followed previous investigators, notably Rabello-Soares, Korzennik, and Schou (2008) and Rhodes et al. (2011). In our case, the simplest criterion is was maximizing mode coverage. Another factor that we consider is considered was choosing intervals similar to each other in length in order to facilitate comparing them.

For the year 2000, only 45 days of continuous data were available, and for 2003, only 38 days were available. There were, however, small additional sections of continuous data for these these years, separated from the previously used time intervals by sections with a low duty cycle. We therefore extended both timeseries time series. In 2002, the situation was reversed; more data were available on the other side of a large gap, but including it did not result in substantially increased mode coverage. Therefore we chose a length that was closer to the other dynamics runs.

The first part of Table 1 shows the timeseries time series that we used for the analysis presented here. The second part of the table shows timeseries time series used in various other investigations. In both cases, processing was carried out through the mode fitting. The timeseries time series and resulting mode parameters can be downloaded from Stanford's the Stanford Joint Science Operation Center (see the Appendix for details). The exception is the 12-day long interval in 2003, which was too short for the mode fitting to succeed, so only timeseries time series are available.

In order to make comparisons with the vw\_V data, we use\_used the same 15 time intervals for two other analyses. Firstly, we use First, we used the regular vw\_V data. Secondly, we use Second, we used the full-disk imagesbut apodize, but apodized them in the same way as the vw\_V data. We also attempted to use the full-disk apodization on vw\_V images that we reconstructed from the full-disk images, but this was only possible for the years 1996 and 1998, because for the other yearsthe gaussian, the Gaussian convolution kernel used for the smoothing reached outside beyond the full-disk crop radius, resulting in the loss of large amounts of data. These last two variations in the analysis required the computation of new leakage matrices. Details of the apodization are provided in the next section.

Table 1. Dynamics timeseries time series. Day numbers refer to the first day of the timeseries time series and are given relative to the MDI epoch of 1 January 1993 00:00:00\_TAI. All timeseries time series begin on the first minute of the start date and end on the last minute of the end date. Duty cycles are given for the raw timeseries time series (DC1) and the timeseries time series after gapfilling (DC2). The number of modes fitted with six 6 a-coefficients (NM6) and with 36 a-coefficients (NM36) is also given. The first part of the table shows the timeseries time series used for this article; the second part shows timeseries time series used for various other investigations.

Day	Length [Days]	Start Date	End Date	DC1	DC2	NM6	NM36
1238	63	23 May 1996	24 Jul 1996	0.93	0.98	2039	1729
1563	93	13 Apr 1997	14 Jul 1997	0.91	0.98	2106	1840
1834	92	09 Jan 1998	10 Apr 1998	0.90	0.97	2132	1862
2262	77	$13~\mathrm{Mar}~1999$	$28~\mathrm{May}~1999$	0.92	0.97	2101	1809
2703	98	$27~\mathrm{May}~2000$	$01~{\rm Sep}~2000$	0.74	0.89	2056	1770
2980	90	$28~{\rm Feb}~2001$	$28~\mathrm{May}~2001$	0.91	0.97	2088	1837
3331	109	$14~{\rm Feb}~2002$	$02~\mathrm{Jun}~2002$	0.85	0.96	2092	1839
3904	76	10  Sep  2003	24 Nov 2003	0.58	0.75	1988	1603
4202	65	$04~\mathrm{Jul}~2004$	$06~{\rm Sep}~2004$	0.87	0.96	2062	1741
4558	67	$25~\mathrm{Jun}~2005$	$30~\mathrm{Aug}~2005$	0.92	0.98	2082	1755
4830	62	$24~\mathrm{Mar}~2006$	$24~\mathrm{May}~2006$	0.89	0.98	2073	1723
5454	58	$08~{\rm Dec}~2007$	$03~{\rm Feb}~2008$	0.87	0.98	2032	1687
5540	64	$03~\mathrm{Mar}~2008$	$05~\mathrm{May}~2008$	0.85	0.96	2088	1740
5981	65	18 May 2009	21 Jul 2009	0.75	0.84	2017	1631
6335	67	$07~\mathrm{May}~2010$	12 Jul 2010	0.85	0.93	2031	1704
2703 a	45	27 May 2000	10 Jul 2000	0.93	1.00	1919	1556
$3296^{\ b}$	27	$10~\mathrm{Jan}~2002$	$05~{\rm Feb}~2002$	0.86	0.93	1864	1127
$3331 ^{\rm b}$	98	$14~{\rm Feb}~2002$	$22~\mathrm{May}~2002$	0.86	0.97	2081	1821
$3368~^{\rm a}$	72	$23~\mathrm{Mar}~2002$	$02~\mathrm{Jun}~2002$	0.90	0.97	2056	1717
$3904^{\ b}$	12	$10~{\rm Sep}~2003$	$21~{\rm Sep}~2003$	0.81	0.98	0	0
$3942~^{\rm a}$	38	18 Oct 2003	24 Nov 2003	0.81	0.94	1921	1367

<sup>&</sup>lt;sup>a</sup> Rabello-Soares, Korzennik, and Schou (2008)

In all cases, we <u>use\_used</u> a window function common<sup>1</sup> to all analyses for each time interval as input to the gapfilling. The result was mainly to discard a large amount of the regular vw\_V data. We did not repeat the analysis of the regular full-disk data using the common window function, but the native window function included at most 0.23 % more data.

The HMI began producing regular science data on 30 April 2010. Since that time, we have been performing medium- $\ell$  analysis of it-the data using 72-day long timeseries time-series in phase with the original MDI medium- $\ell$  timeseries time

<sup>&</sup>lt;sup>b</sup> Rhodes, private communication (2017)

<sup>&</sup>lt;sup>1</sup>The window function is a timeseries time series of ones and zeros designating good and bad data points, respectively. The common window function is the product of two or more other corresponding window functions.

**Table 2.** HMI timeseries time series. Day numbers refer to the first day of the timeseries time series and are given relative to the MDI epoch. Duty cycles are given for the raw timeseries time series (DC1) and the timeseries time series after gapfilling (DC2).

Day	Start Date	DC1	DC2	Day	Start Date	DC1	DC2
6328	30 Apr 2010	0.996	1.000	7408	14 Apr 2013	0.986	0.991
6400	11 Jul 2010	0.982	0.995	7480	$25~\mathrm{Jun}~2013$	0.990	0.997
6472	$21~{\rm Sep}~2010$	0.968	0.995	7552	$05~{\rm Sep}~2013$	0.967	0.997
6544	$02~{\rm Dec}~2010$	0.989	0.995	7624	16 Nov 2013	0.993	0.997
6616	$12~{\rm Feb}~2011$	0.963	0.991	7696	$27~\mathrm{Jan}~2014$	0.969	0.997
6688	$25~\mathrm{Apr}~2011$	0.997	1.000	7768	$09~\mathrm{Apr}~2014$	0.989	0.995
6760	06  Jul  2011	0.987	0.997	7840	$20~\mathrm{Jun}~2014$	0.991	0.997
6832	$16~{\rm Sep}~2011$	0.966	0.991	7912	$31~\mathrm{Aug}~2014$	0.972	1.000
6904	27 Nov 2011	0.990	0.997	7984	11 Nov 2014	0.992	0.997
6976	$07~{\rm Feb}~2012$	0.966	0.997	8056	$22~\mathrm{Jan}~2015$	0.963	0.991
7048	$19 \; \mathrm{Apr} \; 2012$	0.998	1.000	8128	$04~\mathrm{Apr}~2015$	0.989	0.993
7120	$30~\mathrm{Jun}~2012$	0.990	0.997	8200	$15~\mathrm{Jun}~2015$	0.989	0.997
7192	$10~{\rm Sep}~2012$	0.971	0.997	8272	$26~\mathrm{Aug}~2015$	0.970	0.997
7264	21 Nov 2012	0.993	0.997	8344	$06 \ \mathrm{Nov} \ 2015$	0.978	0.990
7336	$01~{\rm Feb}~2013$	0.972	0.997	8416	01 Jan 2016	0.972	0.997

series. The time intervals for which results are presented here are shown in Table 2. We have also created 360-day long timeseries time-series by concatenating the gapfilled 72-day long timeseries time-series.

# 3. Method

The MDI full-disk data are were processed in almost exactly the same way as the vw\_V data, that is, using the updated methodology method described by LS15. The most notable exception is that for the full-disk data, it is possible to use a larger fraction of the input images; whereas the the vw\_V data are apodized with a cosine in fractional image radius from 0.83 to 0.87, the full-disk data are apodized in the same way from 0.90 to 0.95. It should also be noted that each analysis uses a leakage matrix appropriate to the data used. For the full-disk data, the leakage matrix is was calculated as described by LS15, except that the input images are were not convolved with anything. In particular, we have not accounted for any point spread did not account for any point-spread function, but this is expected to have little effect in the medium- $\ell$  regime.

In summary, all analysis analyses of MDI data presented here are were corrected for various geometric effects during spherical harmonic decomposition: image-scale errors, cubic distortion from the instrument optics, misalignment of the CCD, an error in the inclination of the Sun's rotation axis, and a potential tilt of the CCD. The spherical harmonic timeseries are time-series were then detrended and gapfilled as described by LS15, and Fourier transforms of these

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are—were fit to extract the mode parameters. The fitting, or peakbagging, as it is called, takes—took into account horizontal displacement at the solar surface and the distortion of eigenfunctions by the differential rotation (known as the "Woodard effect") (Woodard, 1989). For the native full-disk analysis, the peakbagging is—was also repeated using asymmetric mode profiles in addition to the normally used symmetric profiles.

For the analysis of HMI data, the input images are were already corrected for optical distortion. Hence, the only geometrical correction applied here is was for the inclination error mentioned in the previous paragraph. After the spherical harmonic decomposition, the HMI data are were processed in almost exactly the same way as the MDI full-disk data. In particular, the images are were apodized in the same way, and therefore an identical leakage matrix is was used. The peakbagging is was performed using both symmetric and asymmetric mode profiles for both the 72-day long timeseries and the 360-day long timeseries the 72-day-long and the 360-day-long time series.

In addition, we have created a proxy for the MDI vw\_V data from the HMI data. This is—was done by binning the HMI data by a factor of four to simulate the MDI full-disk data, convolving them with a gaussian Gaussian and retaining only every fifth point in each direction, as described by LS15. The resulting images are—were then apodized in the same way as the MDI vw\_V data, and the peakbagging likewise uses—used the same leakage matrix. We have fit these data only as 72-day long timeseries 72-day-long time series and only using symmetric profiles.

Whether we use used the HMI images in their native resolution or by way of the proxy, the most significant difference with the MDI processing is was in the detrending. Whereas the MDI data needed to have discontinuities in the timeseries manually identified time series identified manually, for the HMI, for HMI—this information can be derived from keywords in the input data. Furthermore, the quality of the HMI data are is more carefully tracked, so that the keywords also provide a reliable measure of what data are expected to be present.

Due—Owing to its orbit and problems with calibration, the HMI spherical harmonic timeseries time-series contain a strong daily oscillation. We therefore detrend detrended them using different parameters than those used for the MDI data. Although in both cases the timeseries are time series were detrended by subtracting Legendre polynomials of degree seven, for HMI these polynomials are the HMI, these polynomials were fit to an interval of 1100 points (825 minutes), which is was advanced by 960 points (720 minutes). In other words, the detrending intervals overlap overlapped by 140 points (105 minutes). For additional details, the reader is referred we refer to LS15.

# 4. Results

#### 4.1. MDI Mode Parameters

In total, we applied four different analyses to all 15 of the dynamics runs. For conciseness, we shall—make use of the following additional labels: fd\_ap90 for the

full-disk analysis using its regular apodization, fd\_ap90\_as for the same thing when fit with asymmetric profiles, and fd\_ap83 for the full-disk data apodized like the vw\_V data. We shall use the label vw\_ap83 when we use the vw\_V with its regular apodization, but note that we processed it using a window function common to all analyses. We shall also use the label vw\_ap90 for the vw\_V data apodized like the full-disk data, but note that this analysis is only available for the 1996 and 1998 dynamics runs. Figure 1 shows the number of modes fitted using six a-coefficients (which parameterize the dependence of the frequencies on m, see LS15) for all of these analyses. As expected from our previous work, the mode coverage for the fd\_ap90\_as analysis as a function of time is basically the same as that for the fd\_ap90 analysis shifted downward. Interestingly, the other two analyses are closer in coverage to the fd\_ap90 analysis, with the exception of the 2003 dynamics run, which by far had the lowest duty cycle. Apparently by far. Apparently, the regular full-disk analysis was less susceptible to this low duty cycle than all of the other analyses. The effect of using asymmetric profiles on the mode parameters themselves will be is discussed in the context of the HMI analysis.

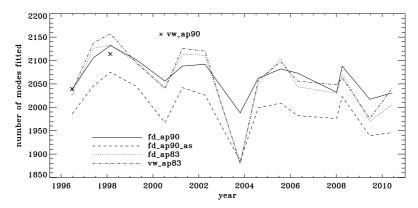


Figure 1. Mode coverage for all dynamics runs. Symbols show the number of modes fitted in 1996 and 1998 by the  $vw\_ap90$  analysis.

In order to compare two different analyses, we must create common modesets. For example, in order to quantify the effect of the apodization, for each dynamics run we find determined the modes common to the fd\_ap90 and fd\_ap83 analyses. For each mode parameter, we then take took a weighted average in time over whatever dynamics runs had each mode successfully fit the dynamics runs that had successfully fit each mode in both analyses. For the weights, we have used the length of each timeseries time series multiplied by its duty cycle. We have also computed the average error, rather than the error on the average, and for comparison between two analyses we use, we used the larger error estimate of the two. Thus the significance that we show is the least that one might expect from an average dynamics run. Lastly Last, the noise parameter [b] requires required special treatment. Since  $e^b$  is proportional to the length of the timeseries time series, each background parameter has had  $\log(T/72.0)$  subtracted from it before

averaging, where T is the length of the timeseries time series in days. Except where noted, we have used the fitted parameters resulting from using six a-coefficients.

In Figure 2 we show the result for six mode parameters: frequency, amplitude, width, background,  $a_1$ , and  $a_2$ . For a full explanation of these, the reader is referred—we refer to LS15. Clearly, the most significant change is to the amplitudes. One might think that this is to be expected since the fd\_ap83 data are apodized to a smaller radius, but in fact, this ought to be corrected for in the leakage matrix. In other words, the parameter A should represent the intrinsic amplitude of the mode on the Sun. Next most significant is the change to the background, which was lower for the fd\_ap83 analysis at lower frequencies, and higher at higher frequencies. The widths were lower for the fd\_ap83 analysis across all frequencies, especially between 2.0 and 3.0 mHz. Lastly Last, although not very significant, the bump seen in the difference in  $a_1$  is encouraging, since it is in the same location as the bump that we hope to eliminate.

Here we note that in the absense absence of systematic errors, these differences should all be small ( $\ll 1\sigma$ ) near the peak power of the p-mode band (around 3.0 mHz), since the signal-to-noise ratio is high there (Libbrecht, 1992). In any case, the differences should have no trends in frequency or any other parameter. One source of random error, the stochastic excitation of the modes, is the same for all observers and apodizations, since the modes considered here have long enough lifetimes lifetimes that are long enough to be considered truly global. Another source of random error, convective motions on the surface, could be different when using different parts of the solar disk, but this still should not cause any offsets in the frequencies, widths, or a-coefficients. Although the amplitudes and background parameters could be affected, such an effect would still be flat in frequency. Even when the signal-to-noise ratio is low, the changes should still be random. Hence, we can already see that there is a problem with the analysis.

To quantify the effect of smoothing and subsampling, we compare the fd\_ap83 and vw\_ap83 analyses in exactly the same fashion. Figure 3 shows the results. Here the convective noise is the same, as well as any instrumental effect, since the two datasets observe almost the same part of the solar disk. Indeed, with the exception of the background parameter, the smoothing and subsampling results in smaller changes than the apodization. The average of the other parameters shows almost no significance at all. In particular, the differences in the a-coefficients are hardly different from zero, which would suggest that the smoothing and subsampling has little effect on any inversion results.

The small difference for the amplitude shown in Figure 3 is, however, deceptive. For all other parameters, the differences look roughly the same for the different dynamics runs, but for the amplitudes, the difference actually alternates in sign. This is shown in Figure 4, where we have plotted the mean significance as a function of time. We have as yet no explanation for this oscillation so far, but focus and tuning changes in the instrument are likely candidates.

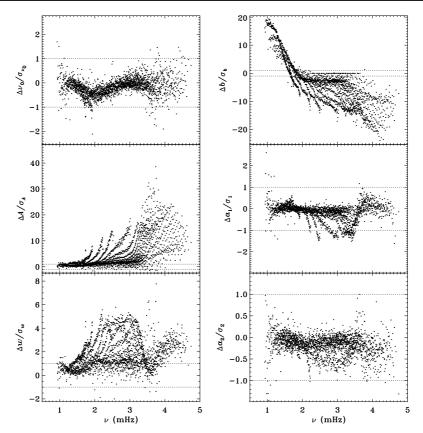


Figure 2. Effect of apodization on mode parameters. Shown are We show changes in frequency  $[\nu_0]$ , amplitude [A], width [w], background parameter [b],  $a_1$ , and  $a_2$  in units of the standard deviation. Each panel is scaled differently; horizontal lines show the  $\pm 1\sigma$  levels. For the a-coefficients, no more than nine points have been excluded from the range shown. The sense of subtraction is fd\_ap90 minus fd\_ap83.

### 4.2. Systematic Errors in MDI data

To explore the effect of the different analyses on our systematic errors, we begin began by performing simple one-dimensional regularized least-squares (RLS) rotational inversions using the  $a_1$ -coefficient onlyalone, just as in LS15. In this case, we formed mode sets common to all three of the fd\_ap90, fd\_ap83, and vw\_ap83 analyses for each dynamics run, and we took the temporal average as before, except that for inversionswe always use, we always used the error on the average. The tradeoff curves in Figure 5 show the result. The curve for the fd\_ap90 analysis has the shape one hopes to see: a single "elbow", so that one may unambiguously choose a tradeoff parameter, not to mention that the  $\chi^2$  values are closer to unity. It is satisfying to see that the value typically used,  $\mu=10^{-6}$ , lies right where it should on the curve: "the place where the residuals stop decreasing sharply, so that further decreases of  $\mu$  will be of little benefit"

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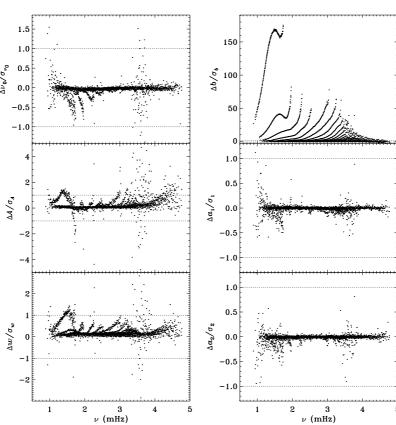


Figure 3. Effect of smoothing and subsampling on mode parameters.

changes in frequency  $[\nu_0]$ , amplitude [A], width [w], background parameter [b],  $a_1$ , and  $a_2$  in units of the standard deviation. For clarity, the bottom panels have at most  $0.65\,\%$  of points excluded. Each panel is scaled differently; horizontal lines show the  $\pm 1\sigma$  levels. The sense of

subtraction is fd\_ap83 minus vw\_ap83.

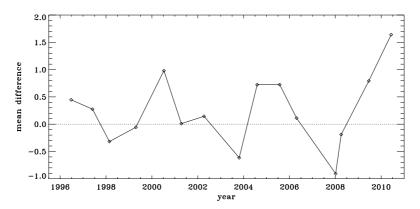


Figure 4. Effect of smoothing and subsampling on amplitudes. Shown are We show the mean changes in units of the standard deviation - for all dynamics runs. The sense of subtraction is fd\_ap83 minus vw\_ap83.

(LS15). The other two curves are very close to the final curve we found in LS15, and we have marked the tradeoff parameters that we used previously.

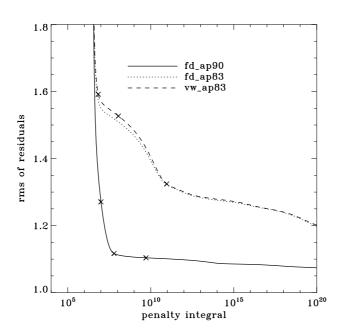


Figure 5. Tradeoff curves for an average over all dynamics runs. Symbols, from left to right, indicate tradeoff parameters of  $\mu=10^{-4}$ ,  $\mu=10^{-6}$ , and  $\mu=10^{-9}$ .

In order to see how the different analyses affect our inference of how the solar rotation varies with latitude, we performed two-dimensional RLS inversions using 36 a-coefficients. First, we form formed averages over the dynamics runs just as we did for the one-dimensional inversions. The residuals of  $a_1$  resulting from inversions of these averages are shown in Figure 6. As one can see, the analyses using the vw\_V apodization clearly show the bump, whereas it is essentially absent from the fd\_ap90 analysis. Investigating the polar jet (LS15), we found that it was clearly visible in inversions of the 1998 dynamics run alone, so we are able to compare all four analyses. Again, we took the modeset mode set common to all four. As Figure 7 shows, we again see that using a smaller apodization radius results in the polar jet, while the larger apodization radius shows no sign of it. Here we must reiterate that the bump does not cause the jet; previous research has shown that excluding modes from the inversion modes that constitute the bump still shows the jet (Schou et al., 2002). Hence, for both the bump and the jet, we are left with a puzzle. Using the vw\_V apodization results in both of the systematic errors, which are then removed by using more data from the input images, although the data added are expected to contain only a small fraction of the helioseismic signal. The most likely explanation is an error in the analysis codes or leakage matrix, but so far, no error explaining our results has been found.

#### 4.3. HMI Mode Parameters

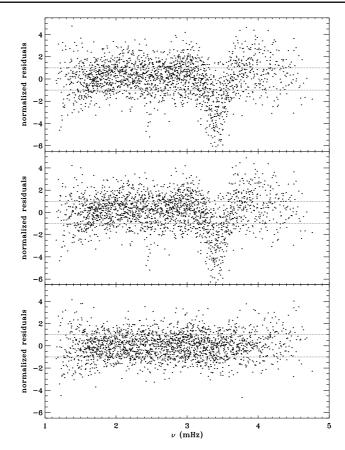


Figure 6. Normalized residuals of  $a_1$  for an average over all dynamics runs. Shown We show from top to bottom are—the vw\_ap83 analysis, the fd\_ap83 analysis, and the fd\_ap90 analysis. Horizontal lines show the  $\pm 1\sigma$  levels.

We have so far So far, we have analyzed about six years of HMI data , as both as 72-day and 360-day fits for the full-disk data, using both symmetric and asymmetric profiles. For the vw\_V proxy, we have used only 72-day long timeseries used only 72-day-long time series and symmetric profiles. The resulting number of modes fitted is shown in Figure 8. The difference in coverage between symmetric and asymmetric fits and between 360-day fits and 72-day fits is what we have come to expect based on our analysis of other datasets. Surprising, however, is the The large oscillation in coverage of the fits to the vw\_V proxy data , is surprising, however, especially since it exceeds the coverage of the full-disk fits at its peak. We shall return to this fact later below.

In LS15, we found that the fits using asymmetric profiles are much less stable than those using symmetric profiles. This is not surprising, since the asymmetric fits require an extra parameter, but it does result in decreased mode coverage. However, in the region where the modes are observed to have strong asymmetry, one must accept that using asymmetric profiles more accurately characterizes them. Hence, the parameters resulting from both types of fitting

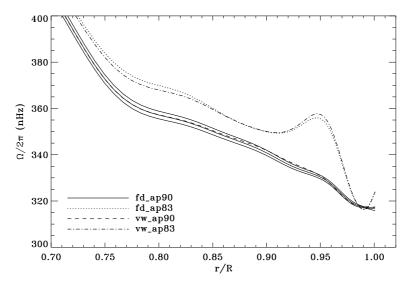


Figure 7. Internal rotation as a function of radius at  $75^{\circ}$  latitude for four analyses applied to the 1998 dynamics run. Solid lines show the fd\_ap90 analysis and its error; errors for the other analyses were similar.

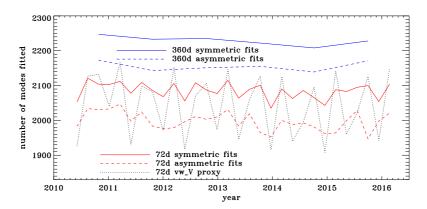
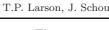


Figure 8. Number of modes fitted as a function of time for the first six years of HMI measurements.

have become standard data products. The difference in coverage for the 72-day fits is shown in Figure 9, where diamonds indicate a mode that failed at least once using symmetric profiles when asymmetric profiles succeeded, and dots indicate a mode that failed at least once using asymmetric profiles when symmetric profiles succeeded. The difference in mode parameters themselves are shown in Figure 10, where we have performed averaging in the same manner as before, using the 72-day fits. This figure is to be compared to the last panel of Figures 4-8 in LS15. Clearly, fitting asymmetric profiles has a large effect on the resulting frequencies in a range between 1.0 and 3.0 mHz. The other mode parameters



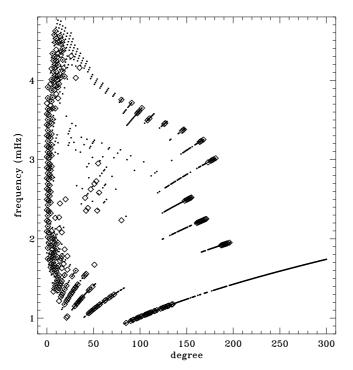


Figure 9. Difference in mode coverage for the first six years of HMI measurements Diamonds show modes that failed to fit at least once with symmetric profiles when asymmetric succeeded, and dots show the opposite.

eters were similarly, but less significantly, affected in a slightly smaller frequency range, still centered at about 2.0 mHz. For the amplitudes, widths, and background parameters, there was also a large and opposite change above 3.8 mHz, while the frequency differences show a second peak around the same frequency. Although not shown here, we found similar differences using the MDI full-disk data. Hence, we can rest assured be sure that the asymmetry of the modes is characterized in the same way by all of the datasets that we have we studied. Unfortunately, this also means that the error magnification that we saw for the frequencies and background parameters in LS15 is also present in the analysis of the full-disk datasets.

Our previous work also revealed discrepancies between 360-day fits and an average over 72-day fits for the MDI vw\_V data, regardless of whether symmetric or asymmetric profiles were used. To confirm that this reflects a characteristic of the algorithm and not the data, we repeated the comparison for the first six years of the HMI. Figure 11 shows the results using asymmetric profiles. Comparison with Figure 13 of LS15 reveals the same trends. The exception is the amplitude differences, but this can be attributed to the gaussian Gaussian smoothing applied to the vw\_V data. Although not shown here, we also found error ratios similar to those shown in LS15. This would indicate that the difference has to do with the algorithm and not with the data. However, Barekat, Schou, and Gizon (2014) found significant differences between the two instruments in the radial gradient of the rotation rate at high latitudes near the surface. In

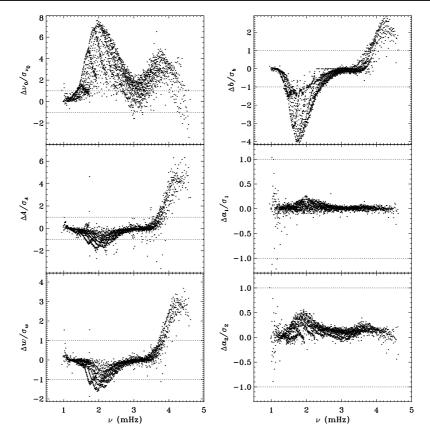


Figure 10. Effect of asymmetric profiles on mode parameters from 72-day fits. Shown are We show changes in frequency  $[\nu_0]$ , amplitude [A], width [w], background [b],  $a_1$ , and  $a_2$  in units of the standard deviation. The data have been averaged over six years of HMI measurements. Each panel is scaled differently; horizontal lines show the  $\pm 1\sigma$  levels. At most, 0.18% of points have been excluded. The sense of subtraction is asymmetric minus symmetric.

subsequent investigations, Barekat (private communication, 2015) also found that the results using the 360-day fits for the HMI differed significantly from those using the averaged 72-day fits, while for the MDI the two are essentially in agreement. Clearly, further study is needed to determine the source of these differences.

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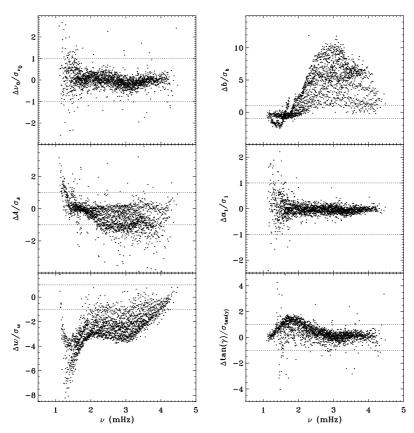


Figure 11. Difference between 360-day and 72-day fits in frequency  $[\nu_0]$ , amplitude [A], width [w], background [b],  $a_1$ , and tangent of the asymmetry parameter  $[\gamma]$  in units of the standard deviation. The data have been averaged over six years of HMI measurements. Each panel is scaled differently; horizontal lines show the  $\pm 1\sigma$  levels. At most, 0.66% of points have been excluded. The sense of subtraction is 360-day long fits  $\underbrace{\text{minus}}_{72\text{-day-long}}$  fits.

#### 4.4. Systematic Errors in HMI data

We plot tradeoff curves and normalized residuals of  $a_1$  for both the HMI full-disk and vw\_V proxy analyses, shown in Figures 12 and 13, just as we did in the same way as for the MDI data. Comparison reveals similar differences between the full-disk and low-resolution results as for MDI. The tradeoff curve shows higher residuals, and the bump in the residuals of  $a_1$  is much more significant. For the rotation rate at high latitudes, HMI's the HMI temporal coverage allowed us to discover that the jet is only discernible when  $|B_0|$  is maximalmaximum, although the two analyses still resulted in significantly different rotation rates. Further Furthermore, the upturn in the rotation rate near the surface at 75° is more pronounced at these times for the vw\_V proxy. When  $B_0$  is close to zero, we see the upturn in both analyses, but it is stronger for the vw\_V proxy. Both features are clearly seen in an average over the six years that we have analyzed, shown in Figure 14.

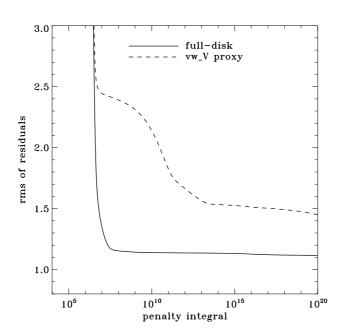


Figure 12. Tradeoff curves for an average over six years of HMI measurements.

## 4.5. Comparison of the MDI and HMI

Both—The MDI and HMI were both operating during the 2010 dynamics run. Hence, we have the opportunity to compare the mode parameters resulting from each dataset. Unfortunately, since the two instruments operate at two different cadences, it is not straightforward to generate a common window function. Setting this aside, Figure 15 shows a comparison of the modes common to the fd\_ap90 and regular HMI analyses for this time interval using their native window functions. Similarly, Figure 16 compares the analysis of the HMI vw\_V proxy and

the MDI vw\_V datasets for the first 72 days of HMI. Again, since realization noise is identical for the two instruments, we hope to  $\frac{\text{to}}{\text{to}}$  see small differences for the frequencies, widths, and a-coefficients, since these parameters should not depend on the  $\frac{\text{height of formation of formation height of the respective absorption lines}}{\text{to the observations}}$ 

These figures are encouraging in that the frequencies and a-coefficients do show little change between the two instruments, although there is a hint of a feature in the frequency differences around 1.7 mHz. One is not surprised to see large differences in amplitude and background parameter, since these parameters do depend on the height at which the mode is observed. The fact that the amplitude differences are not the same in Figures 15 and 16 may be explained by the different center-to-limb dependence of the observing height for the two instruments. Unfortunately, the widths observed by the two instruments are not consistent, with the HMI systematically measuring lower values.

To see how much of the discrepancy results from differences in the instruments and how much results from differences in the processing, Figure 17 plots the difference between the HMI full-disk fits and the fits to the vw\_V proxy data for the first 72 days of HMI measurements, while Figure 18 plots the difference between the fd\_ap90 and vw\_ap83 for the 2010 dynamics run. In other words, Figure 18 can be thought of as the sum of Figures 2 and 3 for a single dynamics run. The close similarity of Figures 17 and 18 gives us confidence that the observed differences have little to do with the source of the data.

## 5. Effect of $B_0$

#### 5.1. Six-Month Periodicity

The original analysis of the vw\_V data revealed a one-year period in the fractional frequency change of the f-modes. In LS15, we found that the amplitude of the annual component increased with increasing degree, but it was decreased by correcting for the Doppler shift that is caused by the motion of SOHO relative to the Sun. In Figure 19 we show the fractional change in f-mode frequency for the entire MDI mission using the most recent fitted mode parameters resulting from using symmetric profiles and 36 a-coefficients. The values shown have been—were averaged over a range in  $\ell$  from 251 to 300 and corrected for Doppler shift. To see how the frequency shifts vary with the solar cycle, we plotted them against the average rms value of the line-of-sight magnetic field, as given by the DATARMS keyword in the corresponding data series<sup>2</sup>. We found a linear relationship between the two and subtracted it. Now, rather than a one-year period, we predominantly see a six-month period, presumably related to

<sup>&</sup>lt;sup>2</sup>For the MDI, this dataseries data series is mdi.fd\_M\_96m\_lev182, which, as the name implies, samples the magnetic field at a cadence of 96 minutes. We looked at all available records for each 72-day interval and rejected outliers above 100 Gauss. The average was taken over the remaining records. For the HMI, the dataseries data series is hmi.M\_720s, which has a cadence of 12 minutes. We therefore took every eighth record to give a sampling similar to that of the MDI, and no outlier rejection was needed.

the absolute value of  $B_0$ . To demonstrate that this is so, we overplot the two quantities in Figure 20. The correlation coefficient between the frequency shifts and the average absolute value of  $B_0$  is 0.42.

To see if the same is true for the HMI, we first apply applied the same procedure to the vw\_V proxy, although in this case, the motion of the spacecraft relative to the Sun has already been corrected for in the dopplergrams Dopplergrams by shifting their target times. To see how the smoothing, subsampling, and apodization might affect the frequency shifts, we did the same with the repeated this for the HMI full-disk data. The result is shown in Figure 21, where we see that the two analyses almost always agree within their errors. In each case, we then subtracted a linear function of the average magnetic field, as before. For the vw\_V proxy, we again see a prominent six-month signal, but it is slightly weaker than for MDI, as Figure 22 shows. In this case, the correlation coefficient was 0.39. For the full-disk data, the correlation was only 0.28. However, inspection of the number of modes fitted as a function of time for the vw\_V proxy, shown in Figure 8, reveals exactly this period. Overplotting the absolute value of  $B_0$  further reveals that -contrary to all expectation, mode coverage is lowest when  $B_0$  is minimal minimum, as Figure 23 shows. Here the correlation coefficient is 0.95. For completeness, we note that the correlation when using the full-disk data is only 0.78. Recalling When we recall that the leakage matrix is was computed assuming  $B_0 = 0$ , it can only come as a shock that we fit more modes when the leakage matrix is most incorrect. Until this discovery, one might have thought that the variation of mode parameters with  $B_0$  was related to the approximation that the leaks from  $\Delta \ell + \Delta m$  odd are zero, since it assumes north-south symmetry. It now seems much more likely that the variation has to do with what which part of the solar surface is visible.

## 5.2. Leaks for Maximum $|B_0|$

A variation in the analysis suggested by the results of the previous section is to use a leakage matrix for a non-zero  $B_0$ . By good fortune,  $B_0$  was near its minimum in the middle of the 1998 dynamics run, its average value being  $-6.35^{\circ}$ . We repeated the peakbagging for this interval using full-disk leakage matrices computed for that this value of  $B_0$  for both apodizations. We must point out, however, that the results using the new leakage matrices are not necessarily any more correct than the original results, since in both cases, the leaks from  $\Delta \ell + \Delta m$  odd are ignored. Put Phrased another way, the leakage-matrix elements that we use become used became more accurate, but the ones that we ignore become those that we ignored became different from zero. To illustrate the relative magnitude of the odd leaks, in Figure 24 we plot sensitivities we plot sensitivities to the target mode ( $\Delta \ell = \Delta m = 0$ ) in the two cases in Figure 24. In Figure 25, we plot odd elements of the new leakage matrices. For the sake of brevity, we have shown clarity, we show only the real part of the radial component of the leakage matrix.

Athough this is not shown, we found that the mode parameters changed similarly for the two apodizations. The unsurprising exception was that the change in  $a_1$  showed the bump, with marginal significance, when using the

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vw\_V apodization. The amplitudes and background parameters showed highly significant changes, while the changes in width were moderately significant. The results of two-dimensional RLS inversions are shown in Figure 26. Clearly, a large change resulted between  $0.83R_{\odot}$  and  $0.95R_{\odot}$  when using the vw\_V apodization was used, whereas the change when using the full-disk apodization was used was not significant. Although this is not shown, we also found similar results using the smoothed data. Plotting the tradeoff curves, shown in Figure 27, we see that the new leakage matrix resulted in lower residuals for both of the apodizations.

# 6. Discussion and Future Prospects

In comparing the MDI full-disk data with the vw\_V data, we found that the difference in mode parameters, with the exception of the background, mostly resulted from the different apodizations used in the two analyses. In particular, the difference in  $a_1$  showed the bump at 3.4 mHz. Correspondingly, two-dimensional RLS inversions of data using the full-disk apodization did not show the bump in the residuals, whereas it appeared almost the same in the two analyses using the vw\_V apodization. Likewise, the high-latitude jet was almost completely absent when using the full-disk apodization. In one-dimensional inversions, the tradeoff curve for the full-disk analysis using the vw\_V apodization still showed the anomalous shape seen in LS15.

To further explore the possible cause of these discrepancies, we plotted the ratio of the amplitudes from the full-disk analysis using its regular apodization to the amplitudes found using the vw\_V apodization, and likewise for the widths. The result is shown in Figure 28. The shape of these ratios is roughly the same as the differences shown in the second and third panels of Figure 2, which were plotted in units of significance. The difference in amplitudes would suggest a problem with the leakage matrix, which could also affect the widths, but those these differences might also be attributed to the model that we use we used for the background. Although not shown, we found that the background differences themselves also showed a trend similar to that seen in the significance.

Smoothing and subsampling made highly significant changes only to the background parameter. Recalling that e<sup>b</sup> multiplies the covariance of the noise at high frequencies (LS15), one might guess that the gaussian Gaussian convolution somehow changes the noise in in that range. The smoothing and subsampling also made significant changes to the amplitude, and these changes varied in sign across the dynamics runs. One probable cause for the sign change is the difference between the best focus and the commanded focus in the instrument, which varied throughout the mission. The occasional changes in the instrument tuning to compensate for drifts are also likely to play a part. The question of how the smoothing and subsampling change the amplitude at all remains unanswered, as their effect should be accounted for in the leakage matrix. In the future, one might perform the smoothing without subsampling, since subsampled data should result in greater interpolation errors when the images are remapped, which could account for some of the differences. Other methods of smoothing

and subsampling are possible, as well as measuring the covariance of the noise in different frequency intervals.

The analysis of HMI data confirmed that using a proxy for the vw-V data resulted in both the high-latitude jet and the bump in the odd a-coefficients, whereas both were essentially absent from the analysis of full-disk data. Comparison of fits using asymmetric mode profiles to those using symmetric profiles revealed differences similar to those seen in LS15 and in the analysis of MDI full-disk data. In spite of fitting fewer modes, asymmetric profiles (occasionally) resulted in more stable fits at the ends of ridges, mostly at the low- $\ell$  ends, but also at the high- $\ell$  ends for p-modes of low to moderate radial order. Comparison of 360-day fits to an average of 72-day fits also revealed differences similar to those seen in LS15. Other investigators (Barekat, Schou, and Gizon, 2016), however, have found differences in the inversions of modesets mode sets from the two instruments, which we have not discussed here, but which should be investigated in the future.

The HMI also allows us to compare the difference between the full-disk results and those for the vw\_V proxy in the magnitude of the six-month oscillation. Although we have not examined did not examine the frequency shifts for the full-disk data, we did find found the surprising result that more modes were fitted for the vw\_V proxy when the absolute value of  $B_0$  was at its peak. This might suggest that the systematic errors we see are related to the alignment of the apodization circles with the spherical harmonic node lines. To see if this is true, one might try using differently shaped apodizations, such as apodizing in longitude and latitude rather than image radius, or an elliptical apodization.

In the comparison of mode parameters from HMI and MDI, we found that differences in frequencies and a-coefficients were not significant for the full-disk analyses, and even less so for the vw\_V analyses. While the frequency differences indicated a small feature, the differences in a-coefficients were almost completely flat. Since these are the only parameters used in rotational inversions, there should be no problem with concatenating datasets from the two instruments in order to increase the interval over which consistent physical inferences can be drawn. As an example, Figures 29 and 30 show internal rotation derived from full-disk datasets for the MDI and HMI, respectively. Following Schou et al. (1998), we have removed the region where estimates of rotation are deemed unreliable. As expected, the two inferences agree quite well.

Furthermore, if we believe assuming that the full-disk analyses are more accurate than the vw\_V analyses, we can use the former to correct the latter. This is essential for the MDI, since the vw\_V data are the only helioseismic dataset it provided with a high duty cycle.

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Disclosure of Potential Conflicts of Interest The authors declare that they have no conflicts of interest.

## **Appendix**

Detailed information on how to access MDI data from the global helioseismology pipeline can be found on the website of the Joint Science Operations Center (JSOC) at jsoc.stanford.edu/MDI/MDI\_Global.html and likewise for the HMI at jsoc.stanford.edu/HMI/Global\_products.html. These pages contain documentation describing how the datasets used in this article were made and how they can be remade. A description of these data and their keywords was also given in the Appendix of LS15; data formats and keyword names remain unchanged in this work. All mode parameters presented here, as well as the rotational inversions shown in Figures 29 and 30, are available in the electronic supplementary material. The data series from which the relevant data may be downloaded are described below.

Mode parameters resulting from both symmetric and asymmetric fits to regular full-disk data from the MDI can be found in two data series: mdi.fd\_V\_sht\_modes and mdi.fd\_V\_sht\_modes\_asym. Mode parameters for the nonstandard analyses can be found in su\_tplarson.mdi\_V\_sht\_modes. In all cases, the first primekey [(T\_START)] should be specified as an MDI day number, found in Table 1, suffixed by "d". Since some of the timeseries time series have the same start times, in general one must also specify NDT, the number of points in the timeseries time series (see Table 1; an MDI timeseries time series has 1440 points per per day). For the nonstandard analyses, the TAG keyword should also be specified; it can take values of fdvwap, vwcomm, and vwfdap corresponding to the labels fd\_ap83, vw\_ap83, and vw\_ap90 used in this article.

For the HMI, all mode parameters presented here reside in the official HMI name space (hmi). The data series are hmi.V\_sht\_modes and hmi.V\_sht\_modes\_asym for the full-disk data, and the day numbers are found in Table 2. For the former data series, there is also a record corresponding to the last dynamics run. Since these series also contain both 72-day and 360-day fits, the primekey NDT should also be specified; note that an HMI timeseries time series has 1920 points per per day. For the vw\_V proxy, the data series is hmi.vw\_V\_sht\_modes.

Newly available online are data series containing the results of two-dimensional RLS inversions for rotation. These series have the same primekeys as those containing the mode parameters and three more in addition: NACOEFF, RADEXP, and LATEXP. NACOEFF is the number of a-coefficients used in fitting the mode parameters, RADEXP is the exponent of the radial tradeoff parameter (=10^RADEXP), and LATEXP is likewise the exponent of the latitudinal tradeoff parameter. To date, only values of RADEXP=-6 and LATEXP=-2 have been used, and these are also the default values for these keywords. NACOEFF can take values of 6, 18, or the default of 36. The data available include the rotation profile, its errors, and the output a-coefficients. The data series names are the same as those given for mode parameters above, with the string modes replaced

by 2drls. For a full explanation of the format of these data<del>the reader is referred</del>, we refer to the electronic supplementary material or the above websites. Note that at this time, all online inversions have used full modesets.

For the mode parameters and inversions in the MDI and HMI namespaces used in this article, the VERSION keyword is set to version2. Furthermore, all of the intermediate data products are also available (archived), and the data series names can be found on the above websites. For the nonstandard analyses, the gapfilled timeseries time series and window functions are archived; the data series are su\_tplarson.mdi\_XXX\_V\_sht\_gf and su\_tplarson.mdi\_XXX\_V\_sht\_gf\_gaps, where XXX can be one of fdvwap, vwcomm, and vwfdap as above. The raw timeseries time-series and window functions have not been archived, but can be recreated if needed.

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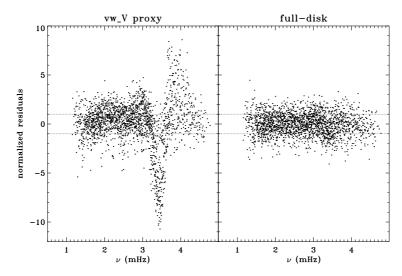
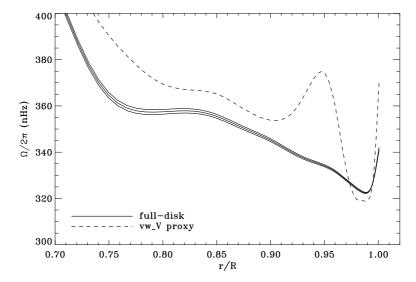


Figure 13. Normalized residuals of  $a_1$  for an average over six years of HMI measurements. Left The left panel shows the HMI vw-V proxy. Right The right panel shows the full-disk HMI analysis. Horizontal lines show the  $\pm 1\sigma$  levels.



**Figure 14.** Internal rotation as a function of radius at 75° latitude for an average over six years of HMI measurements. Solid—The solid lines show the full-disk analysis and its error; errors for the other analysis were similar.

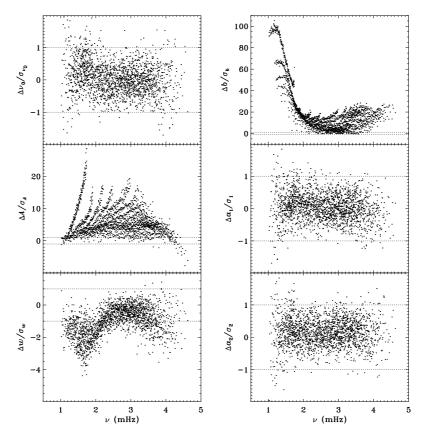


Figure 15. Difference between HMI and MDI full-disk fits for the 2010 dynamics run. Each panel is scaled differently; horizontal lines show the  $\pm 1\sigma$  levels. The sense of subtraction is HMI minus MDI.

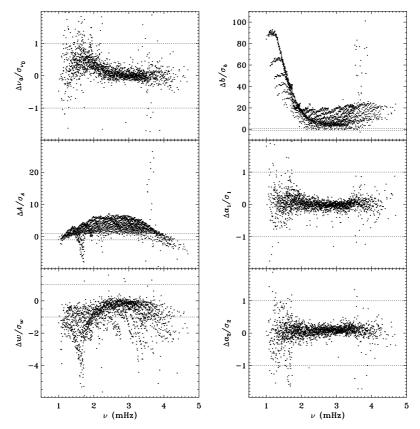


Figure 16. Difference between fits to the HMI vw\_V proxy and MDI vw\_V data for the first 72 days of HMI measurements. Panels are scaled as in Figure 15 for ease of to facilitate comparison, with horizontal lines showing the  $\pm 1\sigma$  levels. The sense of subtraction is HMI minus MDI.

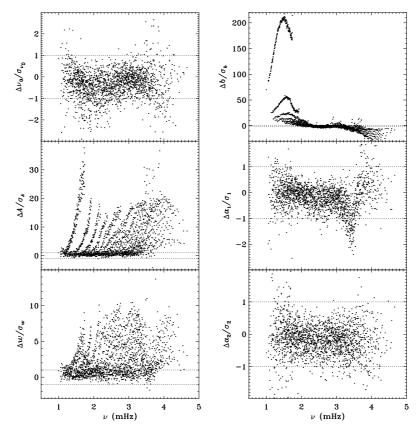


Figure 17. Difference between HMI full disk and vw-V proxy analyses for the first 72 days of HMI measurements. Each panel is scaled differently; horizontal lines show the  $\pm 1\sigma$  levels. The sense of subtraction is full-disk minus vw-V proxy.

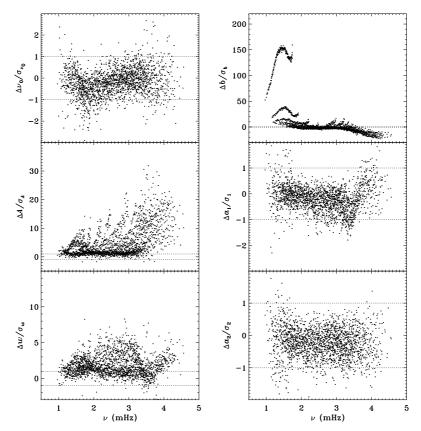


Figure 18. Difference between the MDI fd\_ap90 and vw\_ap83 analyses for the 2010 dynamics run. Panels are scaled as in Figure 17 for ease of to facilitate comparison, with horizontal lines showing the  $\pm 1\sigma$  levels. The sense of subtraction is fd\_ap90 minus vw\_ap83.

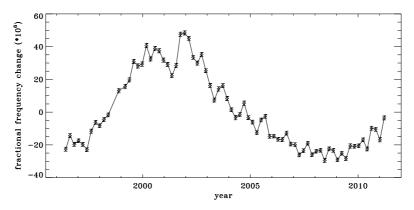


Figure 19. Fractional change in f-mode frequency for the entire MDI mission.

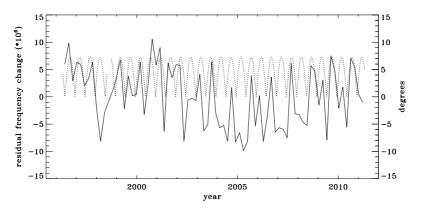


Figure 20. Fractional change in f-mode frequency for the MDI with without the solar-cycle dependence removed. Overplotted is the The absolute value of  $B_0$  is overplotted.

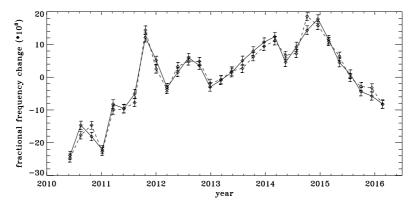


Figure 21. Fractional change in f-mode frequency for the first six years of HMI measurements. Solid The solid line shows the vw\_V proxy, and the dashed line shows full-disk data.

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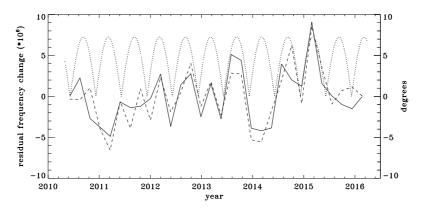


Figure 22. Fractional change in f-mode frequency for the HMI without the solar-cycle dependence removed. Solid The solid line shows the vw\_V proxy, and the dashed line shows full-disk data. Overplotted is the The absolute value of  $B_0$  is overplotted.

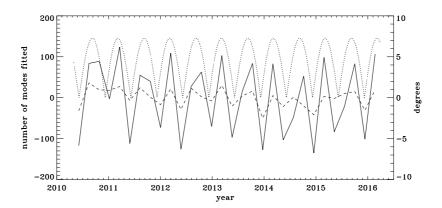


Figure 23. Number of modes fitted as a function of time for the HMI full-disk (dashed line) and the vw\_V proxy (solid line) relative to their means (total number fitted shown in Figure 8). Overplotted is the The absolute value of  $B_0$  is overplotted.

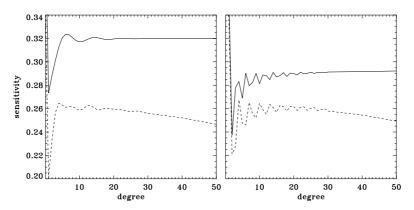


Figure 24. Sensitivity to target mode; the left panel shows m = 0, and the right panel shows  $m = \ell$ . Selid The solid lines show the original leakage matrix, and the dashed lines show leaks for high  $|B_0|$  (= 6.35°).

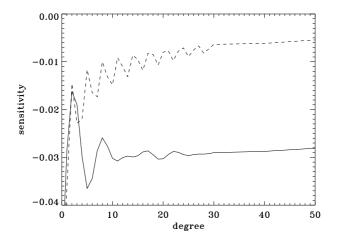


Figure 25. Leaks for  $\Delta \ell = 1$ ,  $\Delta m = 0$  and high  $|B_0|$  (= 6.35°). Solid The solid line shows m = 0, and the dashed line shows  $m = \ell$ . Original The original leaks are identically zero.

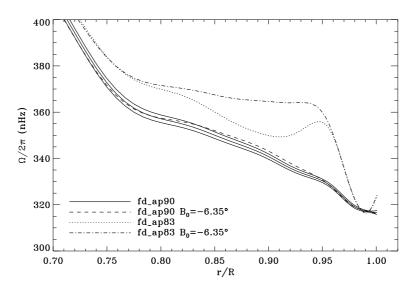


Figure 26. Effect of leakage matrix on inversions. Shown is We show internal rotation as a function of radius at 75° latitude for four analyses applied to the 1998 dynamics run. Two of the curves were shown in Figure 7. Solid—The solid lines show the fd\_ap90 analysis and its error; errors for the other analyses were similar. For these inversions the full modesets, the full modesets were used, rather than common modesets mode-sets.

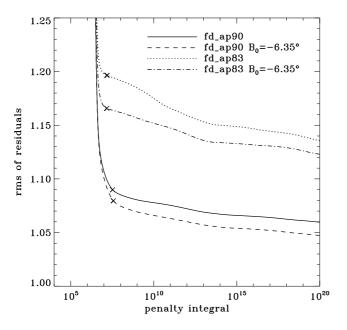


Figure 27. Effect of leakage matrix on residuals. Shown are We show tradeoff curves for four analyses applied to the 1998 dynamics run. Symbols indicate a tradeoff parameter of  $\mu=10^{-6}$ . For these inversions the full mode-sets were used, rather than common modesets mode-sets.

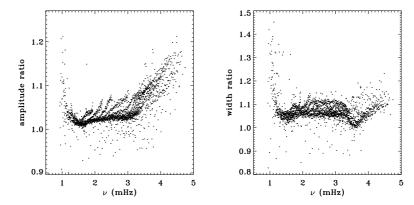


Figure 28. Ratios of amplitude and width from the fd\_ap90 analysis to those from the fd\_ap83 analysis for an average over all dynamics runs. For the width, 17 points have been excluded from the range shown.

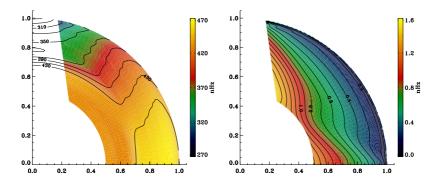


Figure 29. Internal rotation (left) and the corresponding errors (right) derived from the MDI full-disk analysis averaged over all dynamics runs. We have erased color from the regions where estimates of rotation are deemed unreliable; contours are retained on the left for ease of labeling.

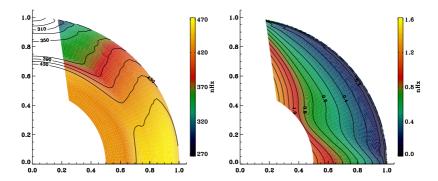


Figure 30. Internal rotation (left) and the corresponding errors (right) derived from an average over the first six years of the HMI 72-day analysis. We have erased color from the regions where estimates of rotation are deemed unreliable; contours are retained on the left for ease of labeling.